

Nested Leaky-Box Model

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Primary Nuclei solution

- ▶ The equilibrium density N_i of a primary cosmic ray species i in the Galaxy is given by

$$n_i = Q_i \tau_G \quad (1)$$

where τ_G is the **energy-independent** galactic escape time and Q_i is the galactic emissivity. We assume NO inelastic interactions of primaries in the cocoon.

- ▶ The emissivity as a function of the kinetic energy per nucleon E is given by

$$Q_i = \rho_s q_i(E) \quad (2)$$

where $\rho_s = \mathcal{R}/V_G$ is the source rate density and q_i is the injection spectrum.

- ▶ The injected spectrum can be parametrized as

$$q_i(E) = q_{0,i} \left(\frac{E}{E_0} \right)^{-\alpha_i} \left[1 + \left(\frac{E}{E_{b,i}} \right)^{\frac{\Delta\alpha}{s}} \right]^s \quad (3)$$

and normalized such that

$$\int_{E_0}^{\infty} dE A E q_p(E) = \xi E_{CR} \rightarrow q_{0,i} \simeq \frac{\xi_i E_{CR} (\alpha_i - 2)}{A E_0^2} \quad (4)$$

- ▶ the equilibrium density can be thereby written

$$N_i = \frac{\mathcal{R} \tau_G}{V_G} q_i(E) = \frac{\xi E_{SN} N(E) \mathcal{R}}{2\pi R_d^2} \frac{\tau_G}{h} \quad (5)$$

Secondary nuclei solution

- ▶ Here we distinguish two contributions, first the standard Galactic component

$$Q_j^G(E) = N_i(E)n_G c\sigma_{i \rightarrow j} \quad (6)$$

where N_i is the equilibrium density of a primary species i .

- ▶ the second contribution is given by the cocoons, here the source term is given by

$$Q_j^C(E) = \rho_s q_i(E)\tau_c(E)n_c c\sigma = N_i(E)n_c c\sigma_{i \rightarrow j} \frac{\tau_c(E)}{\tau_G} \quad (7)$$

as it depends on the primary equilibrium **in the cocoon**.

- ▶ The secondary equilibrium spectrum

$$N_j = (Q_j^G + Q_j^C)\tau_G \quad (8)$$

- ▶ From that follows

$$\frac{N_j}{N_i} = c\sigma_{i \rightarrow j} n_G \left[\tau_G + \frac{n_c}{n_G} \tau_c(E) \right] \quad (9)$$

- ▶ Or in terms of grammage $\chi_i^G = m_p c n_G \tau_G$

$$\frac{N_j}{N_i} = \frac{\sigma_{i \rightarrow j}}{m_p} [\chi_G + \chi_c(E)] \quad (10)$$

- ▶ It can be easily generalized to the case of multiple primaries, in case of two:

$$\frac{N_j}{N_i} = \left(\frac{\sigma_{i \rightarrow j}}{m_p} + \frac{N_k}{N_i} \frac{\sigma_{k \rightarrow j}}{m_p} \right) [\chi_G + \chi_c(E)] \quad (11)$$

- ▶ Assume O purely primary, while C has a primary component plus a secondary one.
- ▶ Oxygen is

$$N_O = \frac{\mathcal{R}\tau_G}{V_G} q_O(E) \quad (12)$$

- ▶ Carbon is

$$N_C = \frac{\mathcal{R}\tau_G}{V_G} q_C(E) + \frac{\sigma_{O \rightarrow C}}{m_p} [\chi_G + \chi_C(E)] N_O \quad (13)$$

- ▶ The ratio is

$$\frac{N_C}{N_O} = \frac{q_C}{q_O} + \frac{\sigma_{O \rightarrow C}}{m_p} [\chi_G + \chi_C(E)] \quad (14)$$

Model parameters

- As in Cowsik et al. the shape of τ_C is taken ad mentulam canis as

$$\chi_C(E) = \chi_C^0 E^{-\zeta \ln E} \quad (15)$$

- Fixed model parameters

$$\mathcal{R} = 1/50 \text{ yr} \quad (16)$$

$$V_G = \pi R_G^2 2h \simeq 2 \times 10^{66} \text{ cm}^3 \quad (17)$$

$$E_0 = 10 \text{ GeV} \quad (18)$$

$$\sigma_{C \rightarrow B} = 60 \text{ mb} \quad (19)$$

$$\sigma_{O \rightarrow B} = 60 \text{ mb} \quad (20)$$

$$\sigma_{O \rightarrow C} = 60 \text{ mb} \quad (21)$$

$$(22)$$

- Free model parameters

$$\chi_G = [1, 10] \text{ gr/cm}^2 \quad (23)$$

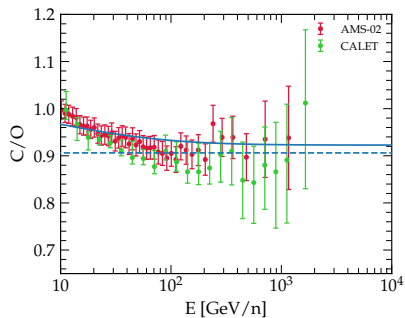
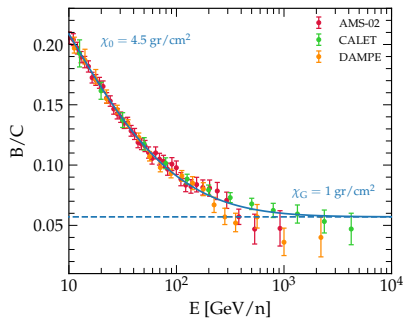
$$\chi_C^0 = [1, 10] \text{ gr/cm}^2 \quad (24)$$

$$\zeta = [0.01, 0.1] \quad (25)$$

$$\frac{N_O}{N_C} = [0.1, 1] \quad (26)$$

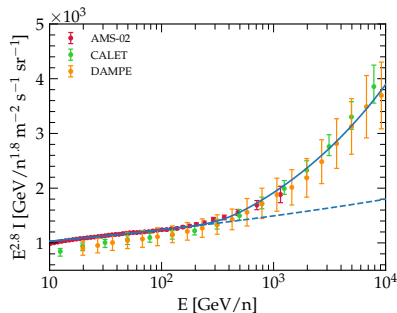
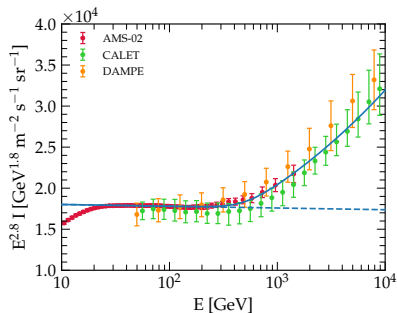
$$(27)$$

Boron-over-carbon



- From the B/C fit $\tau_G \simeq 1 \text{ Myr}$ (for $n_H = 0.5$) which is the energy loss timescale for 1 TeV positrons on CMB only.

Protons



- ▶ From the proton fit $\alpha = 2.81$, $\Delta\alpha = 0.20$, $E_b = 480$ GeV, $\xi = 2\%$
- ▶ From the Helium fit $\alpha = 2.72$, $\Delta\alpha = 0.23$, $E_b = 330$ GeV, $\xi = 0.1\%$

Beryllium

- ▶ Be10 is pure secondary and it decays with decay lifetime τ_d and is given by

$$N_{10} = \left(Q_{10}^G + Q_{10}^C \right) \frac{\tau_G \tau_d}{\tau_G + \tau_d} = c N_i \sigma_{i \rightarrow 10} (n_G \tau_G + n_C \tau_C) \frac{\tau_d}{\tau_G + \tau_d} \quad (28)$$

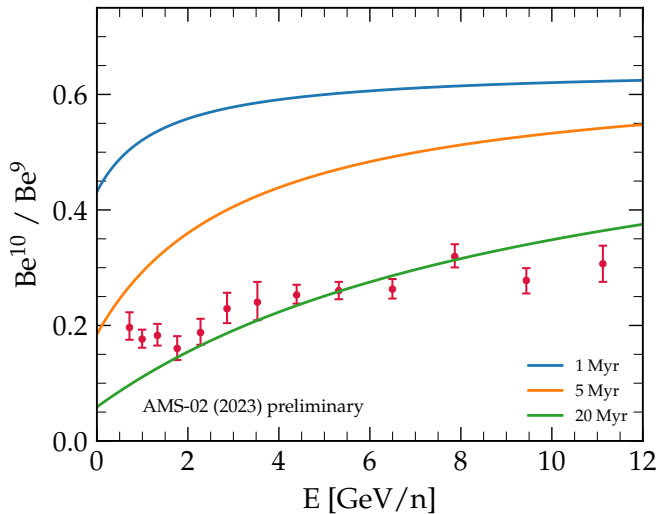
- ▶ following for the ratio

$$\frac{N_{10}}{N_9} = \frac{\sigma_{i \rightarrow 10}}{\sigma_{i \rightarrow 9}} \frac{\tau_G + \tau_d}{\tau_d} \quad (29)$$

- ▶ or in case of more than one primary

$$\frac{N_{10}}{N_9} = \frac{\sigma_{i \rightarrow 10} + \frac{N_k}{N_i} \sigma_{k \rightarrow 10}}{\sigma_{i \rightarrow 9} + \frac{N_k}{N_i} \sigma_{k \rightarrow 9}} \frac{\tau_G + \tau_d}{\tau_d} \quad (30)$$

Beryllium



Antiprotons

- ▶ The antiproton emissivity is again given by the sum of two contributions

$$Q_{\bar{p}} = cn_H \xi \int_{E_t(E)}^{\infty} dE' n_p(E') \frac{d\sigma(E, E')}{dE} \quad (31)$$

- ▶ while in the cocoon is

$$Q_{\bar{p},c}(E) = cn_{H,c} \xi \int_{E_t(E)}^{\infty} dE' \rho_s q_i(E') \tau_c(E') \frac{d\sigma(E, E')}{dE} \quad (32)$$

- ▶ therefore

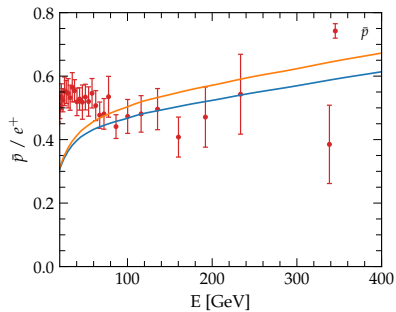
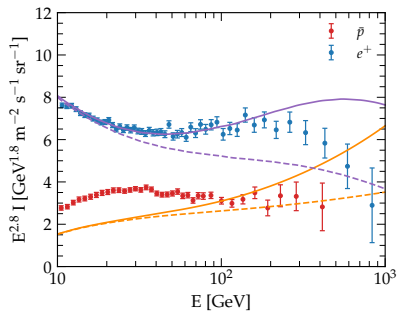
$$N_{\bar{p},G} = Q_{\bar{p},G} \tau_G = \xi \int_{E_t(E)}^{\infty} dE' N_p(E') \frac{\chi_G}{m_p} \frac{d\sigma(E, E')}{dE} \quad (33)$$

$$N_{\bar{p},c} = Q_{\bar{p},c} \tau_G = \xi \int_{E_t(E)}^{\infty} dE' N_p(E') \frac{\chi_c(E')}{m_p} \frac{d\sigma(E, E')}{dE} \quad (34)$$

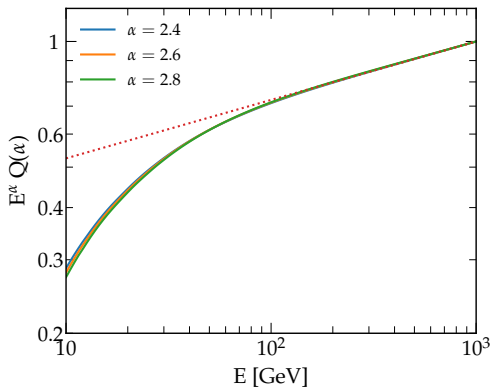
- ▶ For positrons

$$N_{e^+,G} = Q_{e^+,G} \frac{\tau_G \tau_l}{\tau_G + \tau_l} = \xi \frac{\tau_l(E)}{\tau_G + \tau_l(E)} \int_{E_t(E)}^{\infty} dE' N_p(E') \frac{\chi_G}{m_p} \frac{d\sigma(E, E')}{dE} \quad (35)$$

Anti-matter



Anti-matter



$$N_{\bar{p}} \propto E^{-\alpha_{p,\text{He}}} E^{\beta_\sigma} E^{-\delta} \quad (36)$$

$$\frac{N_{\bar{p}}}{N_p} \propto \frac{E^{-\alpha_{p,\text{H}}} E^{\beta_\sigma} E^{-\delta}}{E^{-\alpha_{p,\text{L}}}} = E^{\Delta\alpha} E^{\beta_\sigma} E^{-\delta} \rightarrow 0 = 0.2 + 0.14 - \delta \quad (37)$$